

Design optimization of the plastic scintillators with WLS-fibers and SiPM readouts in the top veto tracker of the JUNO-TAO experiment*

Guang Luo,¹ Y.K.Hor,^{1, †} Peizhi Lu,¹ Zhimin Wang,^{2,3, ‡} Ruhui Li,² Min Li,^{2,3} Yichen Li,^{2,3}
Liang Zhan,^{2,3} Wei Wang,^{1,4, §} Yuehuan Wei,⁴ Yu Chen,¹ Xiang Xiao,¹ and Fengpeng An¹

¹*School of Physics, Sun Yat-sen University, Guangzhou 510275, China*

²*Institute of High Energy Physics, Beijing 100049, China*

³*University of Chinese Academy of Sciences, Beijing 100049, China*

⁴*Sino-French Institute of Nuclear Engineering and Technology, Sun Yat-sen University, Zhuhai 519082, China*

Plastic scintillator (PS) embedding wavelength shifting (WLS) fiber is widely used in high energy particle physics, as muon taggers, and also in medical physics and other applications. In this work, a simulation package is built to evaluate the effects of the diameter and the layout of the optical fiber on the light yield with different configurations. The optimal optical configuration was designed based on the simulation and then validated with two PS prototypes under certain experimental conditions. In the study, the top veto tracker (Tvt) of the JUNO-TAO experiment, comprised of 4 layers of 160 strips of PS was designed and evaluated. When a muon tagging efficiency of a PS strip is higher than 99%, the threshold is evaluated. The efficiency of 3-layer out of 4-layer of Tvt will be higher than 99% even with the tagging efficiency of a single strip as low as 97% using a threshold of 10 p.e. assuming 40% SiPM PDE.

Keywords: Plastic scintillator, WLS-fiber, Light yield, Optical transmission performance, Muon tagging efficiency, JUNO-TAO.

I. INTRODUCTION

The collisions between the primary cosmic rays and the earth's atmosphere will produce a large number of muons[1], the average kinetic energy of which at the sea level is several GeV[2]. Because of the high energy, large mass, small deceleration and deflection in the electromagnetic field, and small bremsstrahlung effect with the atomic nuclear electric field in the matter, the muons will have a strong penetration power[3]. A sub-system of the muon veto detector with high muon tagging efficiency is very important to greatly reduce the background induced by the cosmic-ray (CR) muons for the experiments with only limited overburden near the ground, where the flux of muon is normally four to seven orders of magnitude higher than the underground laboratories with large overburden, such as Jinping underground laboratory[4], Gran Sasso underground laboratory[5] and Canfranc underground laboratory[6]. For example, in neutrino experiments[7–9], dark matter experiments[10–13], neutrino-less double beta decay experiments[14–16], those muon veto systems require muon tagging efficiency higher than 99%. At present, the detectors based on plastic scintillators (PS) have the advantage of easy machining[17–21], flexible structure design, efficient and stable performance[22, 23]. The PS detectors[24], especially with WLS fibers[25–30] and optical photodetectors (Multi-anode PMTs or Silicon PhotoMultipliers(SiPMs)), were used in OPERA[31], MINOS[32], LHAASO[33, 34] and many other experiments[35–38]. Meanwhile, the PS detectors have many applications in geological imaging[39–42], reactor monitoring and other fields[43–45].

The Taishan Antineutrino Observatory (TAO or JUNO-TAO) is a satellite experiment of the Jiangmen Underground Neutrino Observatory (JUNO)[8, 9]. The main purpose of the TAO experiment is to provide a precise neutrino energy reference spectrum for JUNO and benchmark measurements for the nuclear database. TAO detector system will consist of a central detector (CD), an outer shielding and a veto system. The CD will be placed at around 30 m from one core of the Taishan Nuclear Power Plant. The CD consists of a 2.8-ton gadolinium-doped liquid scintillator (LS) filled in a spherical acrylic vessel. The gadolinium-doped LS, as a target material, reacts with neutrinos from the reactor to measure the neutrino energy spectrum[46]. Because TAO only has a limited vertical overburden of 4 m of concrete, the major backgrounds for the TAO experiment are muon spallation products and accidental coincidences, mostly due to the natural radioactivity, the top veto tracker (Tvt) requires to tag muons with an efficiency of higher than 99%.

In this paper, a comparison is realized in sec.II between a simulation based on Geant4[47–49] and a measurement of a prototype of a PS strip with WLS-fiber readout. The light yield results of the experiment and simulation are consistent for passing through muons. In sec.III, the diameter and the layout of the WLS fiber were further checked for higher light yield against with the simulation. An optimized design of the PS strip with WLS-fiber and SiPM readout is proposed for the Tvt system of JUNO-TAO with high light yield and muon tagging efficiency, which will provide a good reference and guidance for the design of PS detector with WLS-fiber. At the same time, the reliability of the optimal design is preliminarily proved by experiments. In sec.IV, with the proposed PS strip design, the expected performance of the TAO Tvt system is demonstrated. Finally, a summary is given in sec.V.

* Supported by School of Physics in Sun Yat-sen University, China.

† Y.K. Hor,heyuanq@mail.sysu.edu.cn

‡ Zhimin Wang,wangzhm@ihep.ac.cn.

§ Wei Wang,wangw223@mail.sysu.edu.cn

63 II. PROTOTYPE OF PS STRIP WITH WLS-FIBER AND 64 SIMULATION

65 Muons will deposit their energy when they pass through
66 and interact with the surrounding materials, and the process
67 of muon energy loss is called muon ionization energy loss[3].
68 The average energy loss per distance (mass thickness) can be
69 described by the Bethe-Blot formula (1) [3, 50]:

$$70 -\frac{dE}{dx} = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \left(\frac{2m_e c^2 \beta^2 \gamma^2 W_{max}}{I^2} \right) - \beta^2 - \frac{\delta}{2} \right], \quad (1)$$

71 where K is a constant, z is the unit charge of the incident
72 muon, m_e and c are the electron mass and the speed of light,
73 respectively. Z and A are the atomic number and mass num-
74 ber of the passing-through matter. W_{Max} is the maximum
75 kinetic energy that can be transferred to an electron when the
76 muon collides with the atom. I is the average excitation en-
77 ergy of the matter. β is the ratio of the speed of a particle to
78 the speed of light. γ is the lorentz factor. δ is the correction
79 factor of the density effect of the matter. The above parame-
80 ters are constant for a given matter.

81 From the formula (1), the deposited energy of muon in
82 the material is related to the energy of muon and the atomic
83 number of the material. For thin-layer media with an atomic
84 number less than 20, such as a PS strip, muons almost pass
85 through in a straight line. Partially lost energy of muon will
86 be converted into light in a PS strip, and the light will be ex-
87 ported by the WLS-fiber, where SiPMs coupled with the fiber
88 is an effective, convenient, and rapid method to pick up the
89 photons and then convert them into an electrical signal.

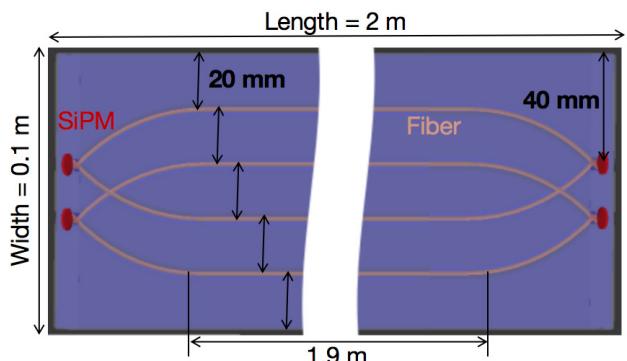


Fig. 1. Design of a prototype of the PS strip with WLS-fiber (Sub-
sequently, it is named option 1)

90 A prototype of the PS with WLS-fiber readout is de-
91 signed and fabricated as shown in Figure 1[51], named op-
92 tion 1. Its size is in a dimension of 2 m (Length) \times 0.1 m
93 (Width) \times 0.02 m (Thickness), and four optical fibers with a
94 diameter of 1 mm are used in total. The pink lines repre-
95 sent the WLS fiber equally spaced, inserted, and filled in the
96 surface of the PS strip. The length of the straight portion is
97 1.9 m. The arrangement is symmetrical in both length and
98 width directions. Two of the fibers are focused into a single

99 group which can be coupled with optical sensors. This op-
100 tion can reduce the number of optical sensors. For example,
101 four SiPMs (red circular point in the figure) can be used for
102 each fiber group, or two PMTs can be used for each PS end.
103 Finally, the PS is wrapped in reflective film (aluminum foil)
104 except for the pips of the optical fiber to export photons, and
105 more details can be found in Ref[51]. The current prototype
106 was tested with CR muon before design optimization.

107 A CR muon survey along the length of the PS strip proto-
108 type with PMT/SiPM is done with the schema shown in Fig-
109 ure 2 (the arrangement of optical fibers is not shown). The CR
110 muon[28, 52] will be selected by the two muon monitors of
111 small scintillators on different hitting locations. When the CR
112 muon monitors are triggered, the signals of the PS strip will
113 be recorded too. 9 points equally spaced along the strip were
114 measured in total relative to the center of the PS. More details
115 about the experiment can be found in Ref [51]. The PS strip
116 is fabricated by Beijing Hoton Nuclear Technology Co. Ltd
117 [53]. The type of fiber is the WLS fiber BCF92 [54, 55]. For
118 further understanding and optimization, a Monte-Carlo (MC)
119 simulation project is set up based on Geant4[47, 49, 52] with
120 the parameters provided by the manufacturer of the PS, the re-
121 flective film and WLS-fiber [55, 56]. The simulation project
122 mainly includes three parts: the first part is the detector ge-
123 ometry, and the PS geometry with optical fibers is designed
124 through a geometric interface. This is why there are vari-
125 ous geometric designs in subsequent optimizations. The sec-
126 ond is the physical process section, which contains a physical
127 list and the optical processes of optical photons. The phys-
128 ical list includes such things as ionization, bremsstrahlung,
129 multiple scattering, pair generation, Compton scattering, and
130 photoelectric effects. Optical processes include the genera-
131 tion of Scintillation and Cherenkov light, wavelength shift
132 effects, Rayleigh scattering, bulk absorption, and boundary
133 processes. The third is the extraction and analysis of informa-
134 tion. The PS and PMT/SiPM are set as sensitive areas; The
135 PS is responsible for obtaining information about the muon,
136 and the PMT/SiPM is responsible for obtaining information
137 about the photons hitting. A parameter interface is provided
138 in the simulation package to set the properties of the mate-
139 rial, such as the PS attenuation length, the scintillator yield
140 (refers to the number of photons generated when the energy
141 deposited in the scintillator is 1 MeV), the reflectivity of re-
142 flective film. By scanning these parameters, a series of simu-
143 lation responses, including the light yield (referring to photo-
144 electron (p.e.) with consideration of the corresponding pho-
145 ton detection efficiency (PDE)/quantum efficiency (QE) of
146 SiPM/PMT) along the PS longitudinal direction, is obtained.
147 χ^2 analysis is performed with the experimental data. Table 1
148 presents a list of parameters corresponding to the minimum
149 χ^2 . In the subsequent optimization work, the parameters of
150 PS, optical fiber and reflective film in table 1 are the same in
151 simulation.

152 The comparison of light yield between the measurement
153 and simulation at different positions is shown in Figure 3(a),
154 where the light yield (in p.e.) refers to the average value of
155 the photoelectron distribution of the selected muons hitting
156 each position and the X-axis represents the distance from the

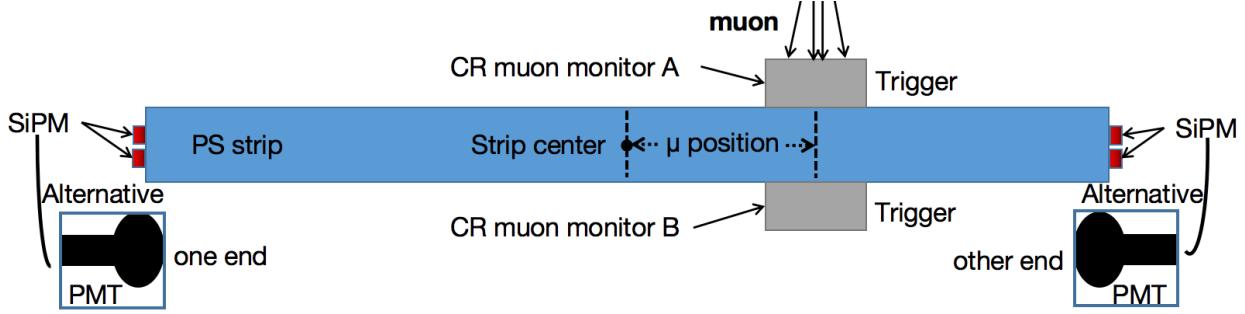


Fig. 2. Device and principle diagram for CR muon measurement with the PS strip prototype. The CR muon will be selected by the two muon monitors of small scintillators on different hitting locations. When the CR muon monitors are triggered, the signals of the PS strip will be recorded too.

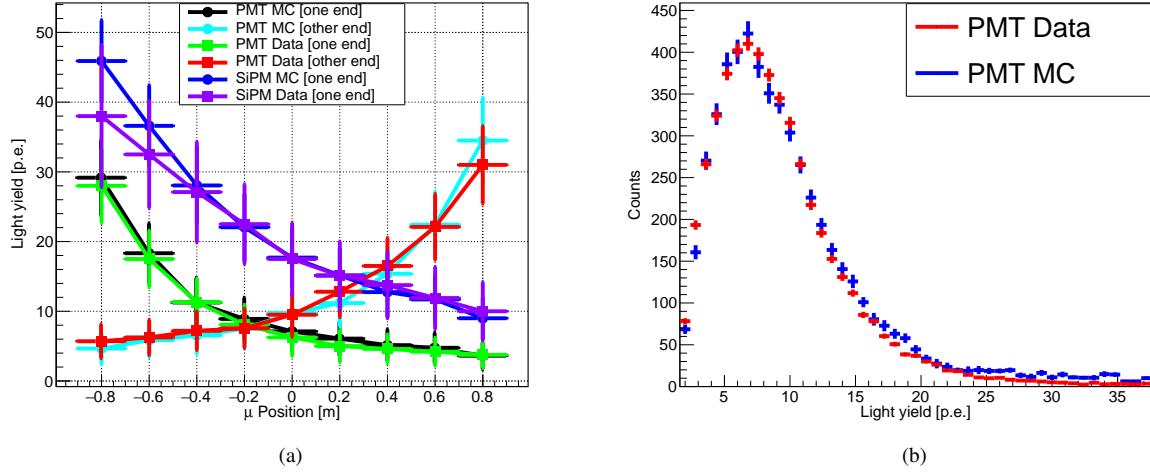


Fig. 3. (a): Comparison of light yield between measurement and simulation at different muon hitting positions. The horizontal error bar is from the dimension of the muon monitors, and the vertical error bar in MC is the statistical error. Due to the different QE of PMTs and coupling, light yield at the two ends is not strictly symmetrical. The light yield of SiPM is larger than that of PMT. Based on the advantages of compact requirements, robustness, and higher light yield than PMT, SiPM is proposed as a sensor for TTV. (b): When the CR muon monitor in the center of PS, a comparison diagram of the experimental energy spectrum and simulation of PMT at one end.

Table 1. The key parameters of the simulation of the PS strip prototype

Material	Properties	Parameters
PS	Base	Polyvinyltoluene
	Scintillation yield (photons/MeV)	8000
	Emission peak (nm)	415
	Attenuation length (cm)	200 @ 400 nm
WLS-Fiber	Core	Polystyrene
	Attenuation length (cm)	380 @ 400 nm
Reflective film	Base	Aluminium
	Reflectivity	85%

center of the PS strip in the length direction. When using the PMT as the sensor, the simulation and data are in good agreement within the error. The light yield at both ends should be symmetrically distributed around the center of the PS. Due

to the different QE of PMT, the light yield at the two ends is not strictly symmetrical. When using SiPM of the same type as a sensor, the distribution of the light yield at one end is shown here. As can be seen from the figure, the light yield of SiPM is larger than that of PMT. Based on the advantages of compact requirements, robustness, and higher light yield than PMT, SiPM is proposed as a sensor for TTV. In addition to the initial comparison shown in Figure 3(a), another similar comparison of energy spectra is also obtained between the measurement and simulation at center of PS shown in Figure 3(b), except for individual energy points that do not conform to the experiment, others are in good agreement with the experiment within the error range. From an experimental perspective, each PMT/SiPM has a different QE/PDE. Since the following optimization work does not pay attention to the impact of electronics, in subsequent studies, the QE of each PMT or the PDE of SiPM will be input with the same value.

III. OPTIMIZATION OF PS STRIP LAYOUT

On the basis of the agreement between experiment and simulation, a further study on the configurations of the PS strip design is done including the diameter and the arrangement of the WLS-fiber for the light yield and tagging efficiency. The relationship between the light yield and fiber diameter is shown in Figure 4 for 2 PS configurations with the same number of fibers embedded. Each point refers to the light yield of the muon hitting the PS center. The magenta heart represents the measurement of the prototype with PMT mentioned earlier. The simulated relationship between the light yield and the fiber diameter was represented by red lines with a PS width of 100 mm and a spacing of approximately 20 mm between neighbor fibers. On the other hand, the blue lines represented the simulated relationship with a PS width of 200 mm and a spacing of around 40 mm between neighbour fibers. Most photons are collected through the WLS fiber by the sensors. Since the area of the optical sensor (Since TAO TVT requirement for readout, subsequent optical sensors default to SiPMs) is larger than the dimension of the optical fiber, there are still some PS scintillation photons directly collected by the sensors without going through the WLS fiber. The two straight lines at the bottom represent the photons directly from the PS scintillation to the SiPM, where the light yield contributed by the PS itself is basically independent of the fiber diameter.

According to the trend of the plot, the larger the fiber diameter, or the smaller spacing of the neighbor fibers, the higher the light yield. Finally, a 1.5 mm fiber diameter is suggested according to the expected response and reasonable cost.

Following the requirements of the JUNO-TAO TVT system, the PS strip with 20 cm width is suggested according to the fabrication, electronics, and cost, but there is still more than one option proposed following a different strategy as shown in Figure 5. The length and thickness of all the options are the same as 2 m, and 20 mm is a basic requirement. The width of option 2, option 3, option 3-1, and option 4 are 200 mm, while the width of option 1 is only 100 mm as the measured prototype. The number of fibers is eight to have a similar spacing except option 1. The fiber of option 2 is only arranged on the front of the PS as doubled of option 1. Option 3 is similar to option 2, but with four optical fibers installed on the back and another four on the front of the PS, respectively. The fiber layout of option 3-1 and option 4 are in a different arrangement. The fiber of options 3-1 is more uniform than that of option 2, but it is not staggered like that of option 4. For option 1, option 2, and option 3, the length of the straight portion of the fiber is 1.9 m. For option 3-1 and option 4, the length of the straight portion of the fiber is 1.5 m.

Table 2. Main configurations of different options of PS strip layout

Configuration	option 1	option 2	option 3	option 3-1	option 4
PS width (mm)	100	200	200	200	200
Fiber numbers	4	8	8	8	8
Fiber diameter (mm)	1	1.5	1.5	1.5	1.5
Fiber spacing (mm)	20	20	20	24	22.5

The PS width and WLS fiber diameter of all the options are listed in table 2. The performance (mainly on the light yield and muon tagging efficiency) and differences among the proposals were further evaluated by simulation. It can be divided into three categories for comparative study.

For Option 2 VS. Option 3: addresses the dependence of light yield on fiber placement.

For Option 2 VS. Option 3-1: addresses the basic dependence of light yield on the uniformity of the fiber arrangement in the PS.

For Option 2 VS. Option 4: addresses the dependence of light yield on the fiber layout.

A. Light yield: $n\text{Photons} \times \text{PDE}$

The angular distribution of the CR muon hitting the PS will affect the PS response. To check the difference and eliminate the additional influence of angular dependence, the muons that are vertically and uniformly incident on the PS strip are used in the simulation firstly. The results of the number of photons received by SiPM for option 1-4 are shown in Figure 6(a). The X-axis is the number of photons ($n\text{Photons}$, light yield = $n\text{Photons} \times \text{PDE}$): the sum of the photons collected by all the SiPMs in a muon event. The figure also shows the average number of photons received by all SiPMs for all muon events. The black line is the distribution of the photons collected by SiPM of option 1. The red and the green lines represent the photon number distributions of option 2 and option 3, which overlap, indicating that the back or front location of fibers has no obvious effect on the collection of the photons. The received photons of the fiber is from the absorption and reemission, reflection of the Cerenkov photons of PS rather than the directly generated photons. The average value of $n\text{Photons}$ in option 2 is almost twice that of option 1. Option 4 has more $n\text{Photons}$ than the previous three options, which indicates that the arrangement of fibers has a large impact on the light yield. The difference of $n\text{Photons}$ between option 4 and option 2 is studied in detail in section III B. All the plots have the same entries, therefore the maximum height of the plots is related to the distribution width.

Another simulation is done to evaluate the differences between option 2, option 3-1, and option 4. The muon generator here is updated in energy and angular distribution according to the Ref.[57, 58] as a more realistic situation to model the response of the PS strip. The results are shown in Figure 6(b). The black, blue, and red lines represent the total photon number distribution of option 2, option 3-1, and option 4, respectively. The figure also shows the average value of photons received by all SiPM of all muon events. Except for the total number of photons, the average value of photons of option 4 is the highest, option 3-1 is the middle, and option 2 is the lowest. Option 4 has the best ability to collect photons.

From the formula (1), it can be known that the deposited energy is directly proportional to the track length of muon

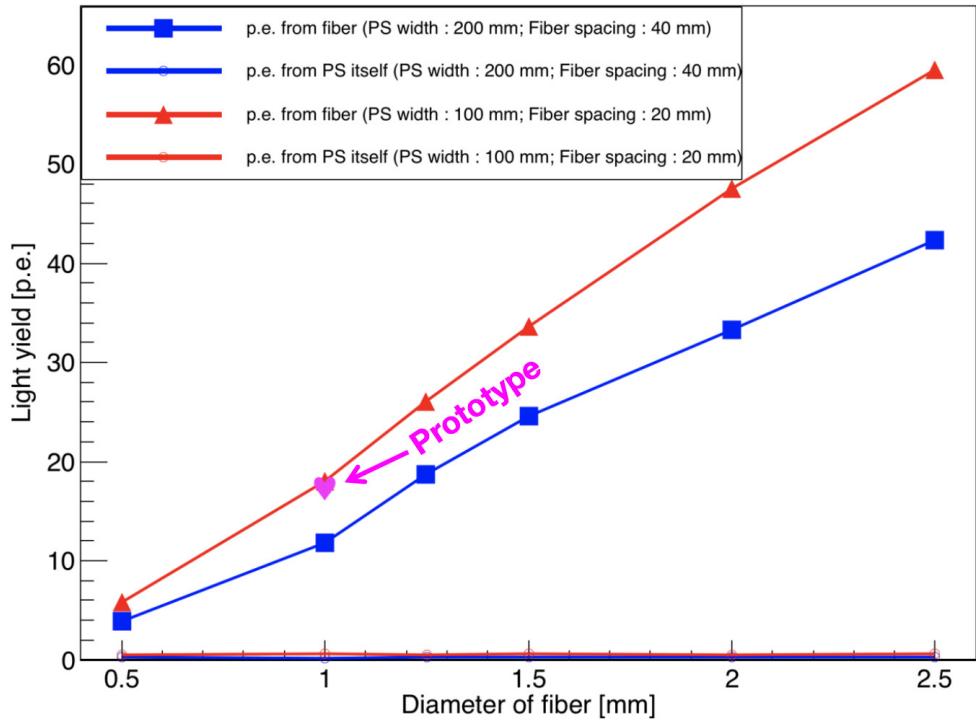


Fig. 4. Assuming 30% SiPM/PMT PDE/QE; each point refers to the light yield of the muon hitting the PS center. The magenta heart represents the measurement of the prototype with PMT mentioned earlier. Two widths of 200 mm and 100 mm for the PS are checked with the same length of 2 m and thickness of 0.02 m. The simulated relationship between the light yield and the fiber diameter was represented by red lines with a spacing of approximately 20 mm between neighbor fibers. The blue lines represented the simulated relationship with a spacing of around 40 mm between neighbour fibers. The same four fibers are used in each case, which is why the light yield of the 200 mm wide PS is smaller than that of the 100 mm one.

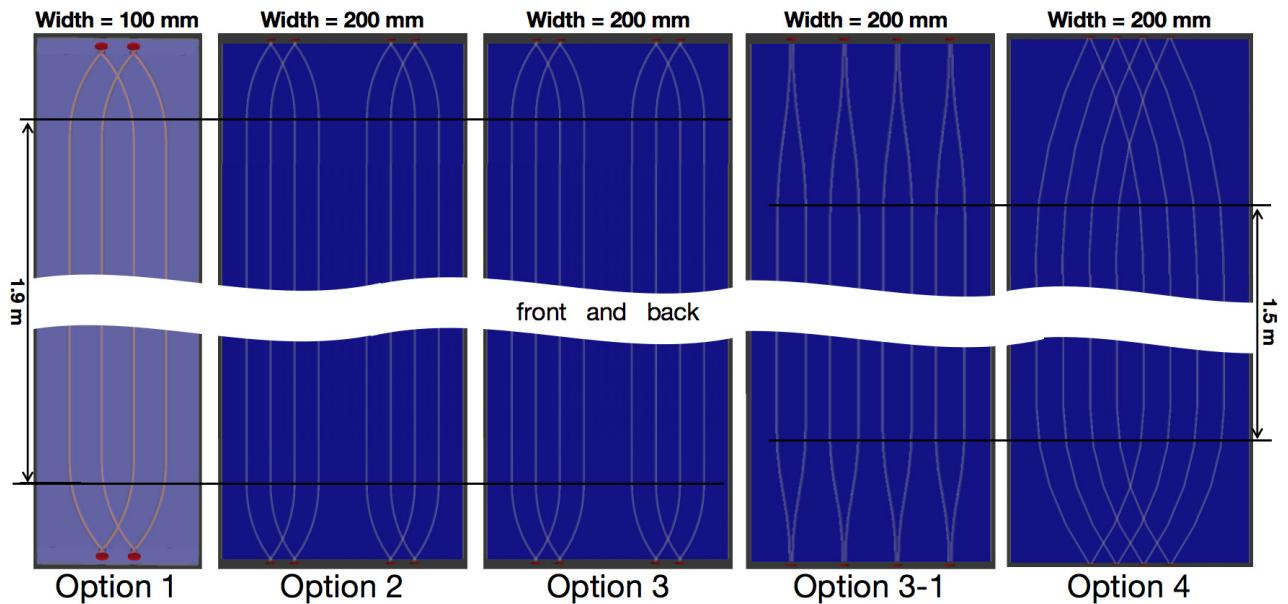


Fig. 5. All evaluated designs of a single PS strip. The number of fibers is eight with a similar spacing except option 1. The fiber of option 2 is only arranged on the front of the PS as doubled of option 1. Option 3 is similar to option 2, but with four optical fibers installed on the back and another four on the front of the PS, respectively. The fiber layout of option 3-1 and option 4 are in a different arrangement. The fiber of options 3-1 is more uniform than that of option 2, but it is not staggered like that of option 4.

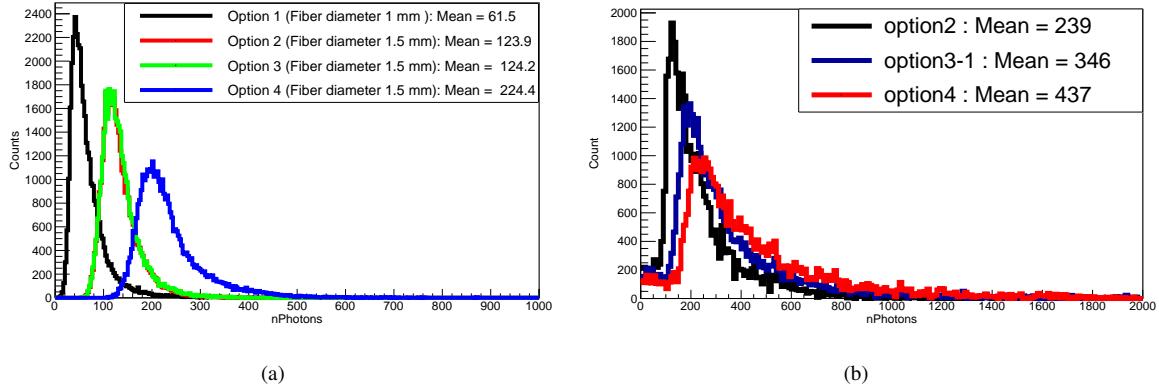


Fig. 6. (a): With vertical incident muon simulation: Distribution of total number of photons received by SiPM for four different options of PS strip layout. nPhotons: the number of photons received/collected by all the SiPMs in a muon event, without the consideration of the corresponding PDE of SiPM. Light yield = nPhotons \times PDE; the average value of nPhotons in option 2 is almost twice that of option 1. Option 4 has more nPhotons than the previous three options. (b): With realistic muon simulation: Distribution of the total number of photons received by SiPM with different fiber arrangements. The black, blue, and red lines represent the total photon number distribution of option 2, option 3-1, and option 4.

chinaXiv:202302.00256v2

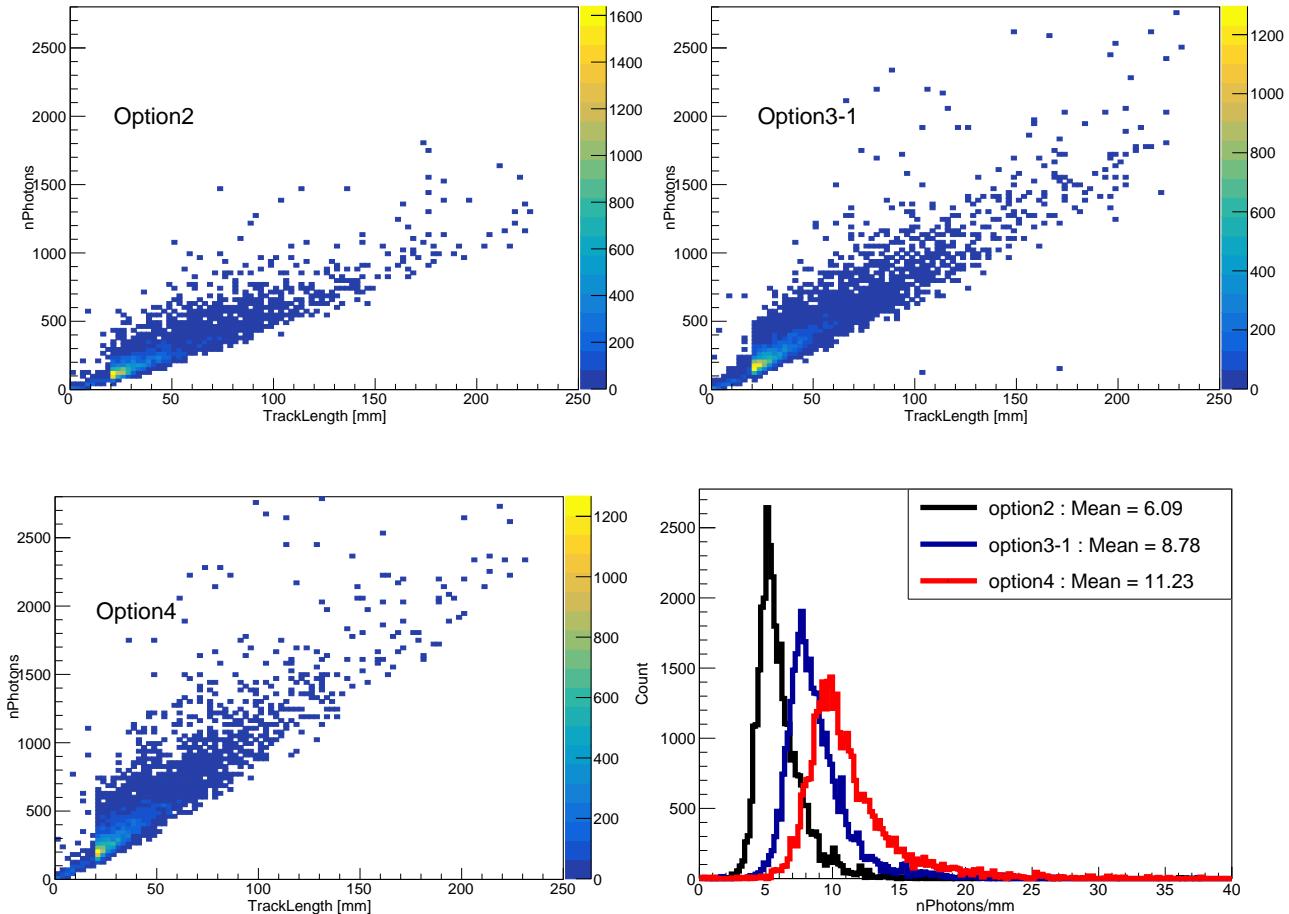


Fig. 7. Total photon number versus muon track length. The upper left, upper right, and bottom right of the picture represent the two-dimensional diagram of total photon number versus muon track length for option 2, option 3-1, and option 4 respectively. Bottom right: comparison of the photon number normalized to unit muon track length.

passing through the material. Because of the oblique incidence of muons of the updated muon generator in the second simulation, the track length of the muons in the PS strip can be either less(represent an edge events) or greater than the thickness of the PS strip. The total photon number in Figure 6(b) is much higher than that in Figure 6(a) even with the same option: the average value of nPhotons has nearly doubled, while some signals with small amplitude show up. The upper left panel of Figure 7 shows the two-dimensional diagram of total photon number versus muon track length for option 2, the upper right panel is for option 3-1, and the bottom left panel is for option 4. The colors of the three panels represent the density of events. The total photon number is proportional to the track length as expected. When the track length is 20 mm (thickness of the PS), the event density is the highest due to the maximum flux density when muon is vertically incident. The bottom right panel of Figure 7 shows the photon number per millimeter of muons passing through the PS strip of the three options. It is obvious that the average value of photons per unit length (mm) of option 4 is 11.23, which is nearly twice that of option 2. This concludes that option 4 is the most effective one for the light yield response to muon when the muon deposits the same amount of energy among these options.

304

B. Transmission performance

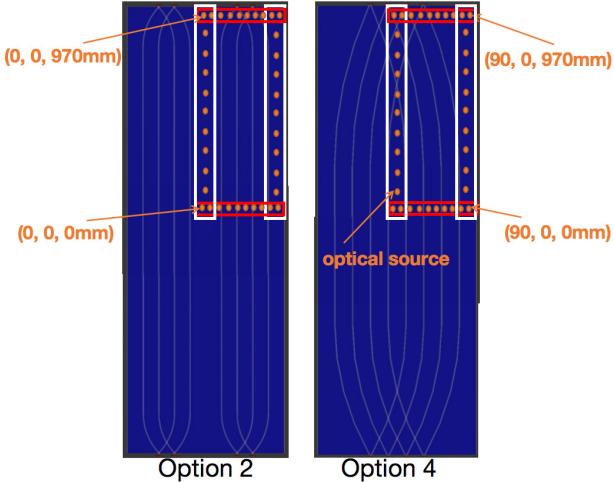


Fig. 8. Position distribution diagram of optical scanning of option 2 and option 4: the two white frames indicate the simulated points of 11 different positions along the length direction of the PS, respectively; and the two red frames indicate the simulated points of 10 different positions along the width direction of the PS, respectively. Each point is equidistant.

Option 2 and option 4 have a similar configuration on WLS-fiber number, PS length, width, and thickness, but their light yields show an obvious difference. The photon transmission process is further checked for a better understanding of any specific reasons for the influence of fiber arrangement

on the light yield for possible further optimization.

Muon deposits its energy along its track in the PS and emits photons in four π directions. Meanwhile, the photons in the PS will propagate through attenuation, absorption, and re-emission effects, so it is not easy to accurately characterize the light transmission performance between different options. A specified optical survey is done to check the photon reflection times before absorption by the WLS-fiber of each photon generated by PS, where the survey locations are shown in Figure 8 trying to cover the center and edge of the PS strip. The orange points are the specified locations to generate optical photons. At each point, 15000 photons are generated in four π directions to mimic the random photons excited by a muon. It is obvious that before the photons enter the fiber, more reflection times, more difficult to reach SiPM (inverse relationship). At the same time, if the number of photons entering the fiber is more, the number of photons arriving at SiPM is more (proportional relationship). In order to model the combination of the two factors, an R-value proposed as the total number of photons entering the fiber divided by the average number of reflections of photons before entering the fiber.

Figure 9 shows the distribution of R-value of an optical survey at different positions under option 2 and option 4. The abscissa of each point is the position corresponding to the dash # symbol in the figure, the specific location can be seen in combination with Figure 8. The ordinate of each point is the R-value stands for optical transmission performance. The red line is the R-value of option 4, and the black represent option 2. The upper left panel of the figure shows the distribution of the R-value along the length direction of the PS strip when the position is the center of the width of the PS. The R-value of option 4 is higher than option 2 in the whole length direction. The lower left panel shows the distribution of the R-value along the length direction of the PS strip when the position in the edge of the width of the PS strip. The R-values of the two are nearly the same, except for the edge in the length direction, where the R-value of option 2 is greater than option 4. The upper right panel reveals the distribution of the R-value along the width direction of the PS when the position is the center of the length of the PS, where option 4 and option 2 have a trend of trade-offs. The lower right panel reveals the distribution of the R-value along the width direction of the PS when the position is the edge of the length direction of the PS, the R-value of option 4 is larger than that of option 2 at the positions within 30 mm from the middle in the width direction of the PS.

The results show that when the photon is generated at a location where the fiber is sparse, the number of photons entering the optical fiber will decrease sharply, while the times of reflections will also increase. Meanwhile, when photons is generated in dense places of fibers, the number of photons entering the fiber will increase dramatically, while the times of reflections will also decrease dramatically.

Looking at the four figures as a global picture, with option 2 configuration, the maximum vertical coordinate of the black point is less than 200, but with option 4 configuration, the ordinate of the red point can be greater than 200. The R-value of option 4 and option 2 has a trend of trade-offs and

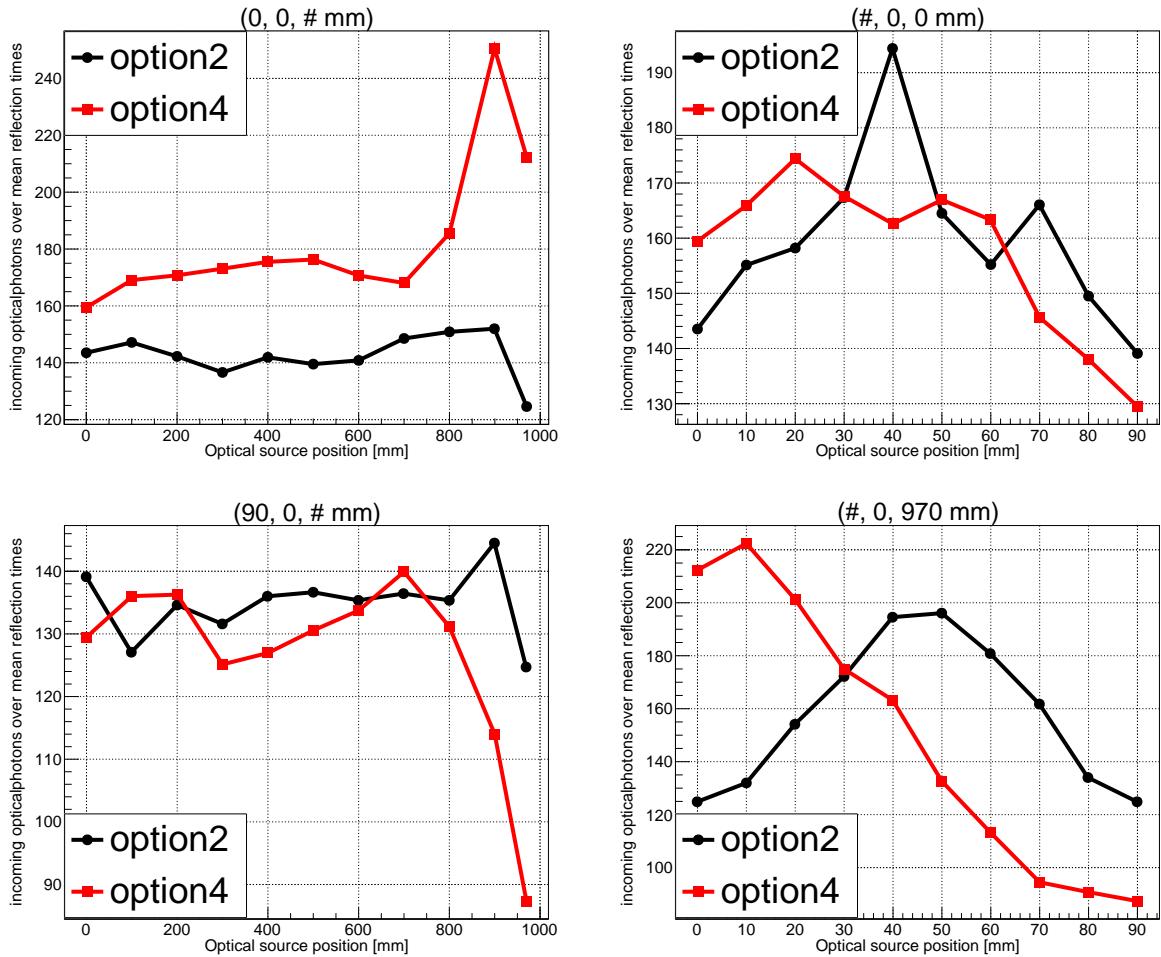


Fig. 9. The distribution of the ratio of the photon numbers entering the fiber to the average times of reflections at different positions of the optical survey. The value of the X-axis of each point is the position corresponding to the dash # symbol in the graph, and the value of the Y-axis of each point in the graph represents the R-value. The upper left panel of the figure shows the distribution of the R-value along the length direction of the PS strip when the position is the center of the width of the PS. The R-value of option 4 is higher than option 2 in the whole length direction. The lower left panel shows the distribution of the R-value along the length direction of the PS strip when the position in the edge of the width of the PS strip. The R-values of the two are nearly the same, except for the edge in the length direction, where the R-value of option 2 is greater than option 4. The upper right panel reveals the distribution of the R-value along the width direction of the PS when the position is the center of the length of the PS, where option 4 and option 2 have a trend of trade-offs. The lower right panel reveals the distribution of the R-value along the width direction of the PS when the position is the edge of the length direction of the PS. The results show that when the photon is generated at a location where the fiber is sparse, the number of photons entering the optical fiber will decrease sharply, while the times of reflections will also increase. Meanwhile, when photons are generated in dense places of fibers, the number of photons entering the fiber will increase dramatically, while the times of reflections will also decrease dramatically.

368 the overall difference is small from the lower left and upper
 369 right panel of Figure 9. The R-value of option 4 and option 2
 370 has a big difference from the upper left and lower right panel
 371 of Figure 9. The distribution of R values in the PS region has
 372 been get. To obtain a quantitative relationship between the
 373 optical transmission performance of option 2 and option 4 in
 374 the entire region, an explanation of mathematical integration
 375 (R-value) were provided. The integral area under the red and
 376 black lines in the upper left and lower right panels of Figure 9
 377 is calculated. Then, the integral area of the red line in the up-
 378 per left panel is multiplied by the area in the lower right panel

379 to obtain the transmission performance of the whole region of
 380 option 4. The same procedure is followed for the black line
 381 as for the red line, and the product of the two integral areas
 382 is used as the transmission performance of the entire region
 383 of option 2. Finally, the product of the integral area of the
 384 red line is divided by the product of the black line, and it is
 385 found that the ratio is 2.21. This shows that the transmission
 386 performance of the whole area of option 4 is 2.21 times better
 387 than that of option 2. Therefore, it can be understood why the
 388 overall photon number of option 4 is more than that of option
 389 2. At the same time, the curve shows that the R-value distri-

390 bution range of option 4 is larger than that of option 2, which
 391 means that the optical export uniformity of option 4 is worse
 392 than that of option 2. This is the result of uneven optical fiber
 393 layout.

394 In short, it was found that if the optical fiber in the PS mid-
 395 dle region is in higher density, the final effective light yield
 396 will be higher. Therefore, to achieve a higher light yield, it is
 397 possible to consider placing as many optical fibers as possible
 398 in the middle of the PS.

399 So far, in terms of light yield, many simulations and made
 400 massive comparisons have done, the best configuration has
 401 been obtained that is option 4 under the current design con-
 402 ditions. At the same time, to some extent, the reason why it
 403 is the best option has been quantitatively explained and good
 404 reference and guidance suggestions for the processing tech-
 405 nology and structure design were provided. For a PS detec-
 406 tor, the most important factor, in addition to light yield, is the
 407 muon tagging efficiency. In the next section, the muon tag-
 408 ging efficiency and inefficiency are discussed.

409 C. Muon tagging efficiency and inefficiency

410 To simplify the simulation process, assuming that the PDE
 411 of SiPM is equal to 30%. The sum of four SiPM outputs at
 412 both ends can be acquired under different options to further
 413 study the muon tagging efficiency. Figure 10 shows the two-
 414 dimensional distribution of light yield at both ends of option
 415 2, option 3-1, and option 4 in simulation, respectively. The
 416 abscissa is the sum of the SiPM outputs at one PS end, and
 417 the ordinate is the sum at the other end. The range of light
 418 yield at both ends of option 4 is the strongest and most di-
 419 vergent. If the sums at both ends of an event are greater than
 420 their respective threshold at the same time, it is considered
 421 that this muon has been detected. The event count of muon
 422 hitting the PS strip is noted as N_{all} , with the event count of
 423 detected muon noted as N_{tag} . The muon tagging efficiency
 424 is defined by the ratio N_{tag}/N_{all} . Here, the environmental
 425 background is not considered, mainly to eliminate environ-
 426 mental background interference. Figure 11 shows the rela-
 427 tionship between the muon tagging efficiency and the thresh-
 428 old of option 2, option 3-1, and option 4, respectively. The
 429 black, blue and red represent the relationship between effi-
 430 ciency and threshold when the PDE of SiPM is 20%, 30%
 431 and 40%, respectively. As seen from the three figures, when
 432 PDE is 20% and the threshold is set to 10 p.e., the tagging
 433 efficiency of option 4 is still higher than 90%, while option
 434 2 and option 3-1 are both less than 90%. With the increase
 435 of the threshold, the efficiency of option 2 decreases fastest,
 436 followed by option 3-1, and option 4 is the slowest. Accord-
 437 ing to Figure 11, the required threshold can be found when a
 438 certain tagging efficiency needs to be achieved.

439 Figure 12 shows the corresponding threshold when the
 440 muon tagging efficiency reaches 99% under different PDEs
 441 of SiPM. When the PDE of SiPM is 40%, for option 2, the
 442 threshold cannot exceed 1.5 p.e.; for option 3, it cannot ex-
 443 ceed 2.4 p.e.; however, for option 4, the threshold can be set
 444 to 3 p.e., which greatly reduces the contribution from dark

445 noise of SiPM. So far, option 4 is superior to other options in
 446 terms of light yield and muon tagging efficiency. Therefore,
 447 Option 4 is taken as the current optimal configuration. When
 448 the PDE of SiPM is 40%, The threshold and corresponding
 449 tagging efficiency of a single PS strip were summarized in
 450 table 3.

Table 3. Threshold and corresponding tagging efficiency of single PS strip with 40% SiPM PDE.

Threshold(p.e.)	3.1	6.3	10	15	19
Tagging efficiency	99%	98%	97%	96%	95%

451 To further verify the simulation, another prototype with op-
 452 tion 4 was built, the parameters of PS, optical fiber and re-
 453 flective film are the same with before prototype in simulation.
 454 Only the PDE of the backend SiPMs is different. 2 SiPM were
 455 used as sensors and using the oscilloscope to collect data. The
 456 SiPM is from the K-series MicroK-40035-TSV[59]. The data
 457 is saved by the lecroy-HDO4104A oscilloscope. Preliminary
 458 experimental results were obtained.

459 Figure 13 shows the experimental measurement and simu-
 460 lation results of the PS prototype of option 4 when the muon
 461 monitor is in the PS center. To avoid the effect of the SiPM
 462 dark noise and environment background. Signals stronger
 463 than 3 p.e. was analyzed, and the blue line is the energy spec-
 464 trum of one SiPM from the measurement. The red line is
 465 the result of the simulation. Every point has an error bar, the
 466 abscissa is the light yield, and the ordinate is the log stand
 467 for the normalized event rate. It can be seen from the figure
 468 that in the spectrum below 12 p.e., the simulated spectrum al-
 469 most corresponds to the experiment. It was found that the
 470 most probable signal amplitude for passing through muons is
 471 around 8 p.e. from data. So for the final configuration used
 472 for TAO where there are 4 SiPMs at each end of the PS strip,
 473 the sum of the most probable signal amplitude will be about
 474 32 p.e. When the spectrum is in the range of 15 p.e. to 25 p.e.,
 475 within the range of error, the event rate of the simulated spec-
 476 trum is more than that of the experiment. When the spectrum
 477 is greater than 25 p.e., and the event rate of simulation and
 478 experiment events is relatively small. In the high energy re-
 479 gion, the energy spectrum does not show good consistency,
 480 but it does not affect our optimization work. Several main
 481 factors lead to differences in the energy spectrum. The first
 482 is that the coupling between SiPM and PS is not perfect and
 483 the same for simulation and measurement. The second is that
 484 the widths of the two small plastic modules for CR monitors
 485 are not the same, which may lead to a deviation of several
 486 centimeters in the location where muon hits. Resulting in an
 487 impact on the energy spectrum.

488 In any case, the light yield of option 4 is indeed much
 489 higher than that of option 1 before optimization. This shows
 490 that our optimization method is effective. According to the
 491 above analysis and design, The light yield and the muon tag-
 492 ging efficiency corresponding to option 2, option 3-1, and op-
 493 tion 4 have been obtained. When the muon tagging efficiency
 494 reaches 99%, option 4 has the highest threshold limit.

495 To study and discuss the inefficiencies by simulation, the
 496 track length of all muons in the PS has been counted. The

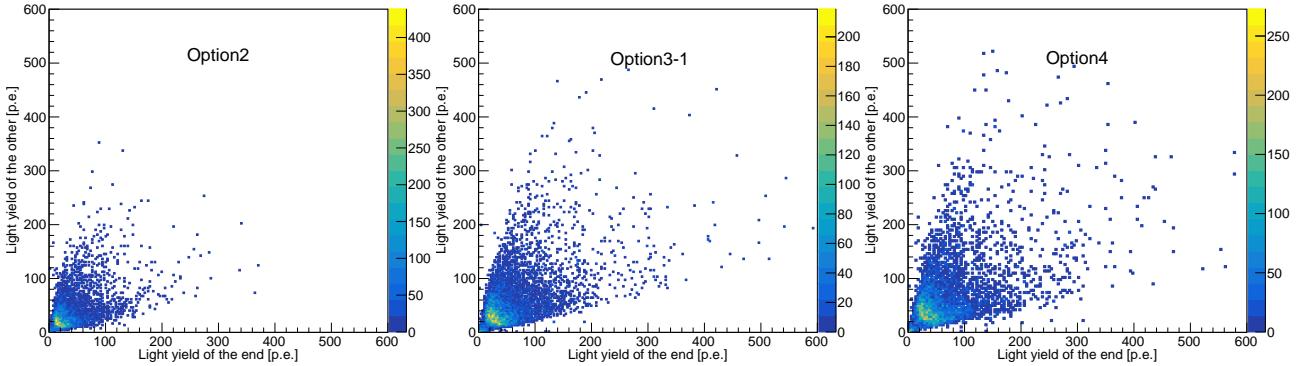


Fig. 10. Assuming 30% SiPM PDE. From left to right: two-dimensional distribution of light yield at both ends of option 2, option 3-1 and option 4, respectively. The color represents the density of events. From options 2 to 4, the distribution range of light yield gradually widens.

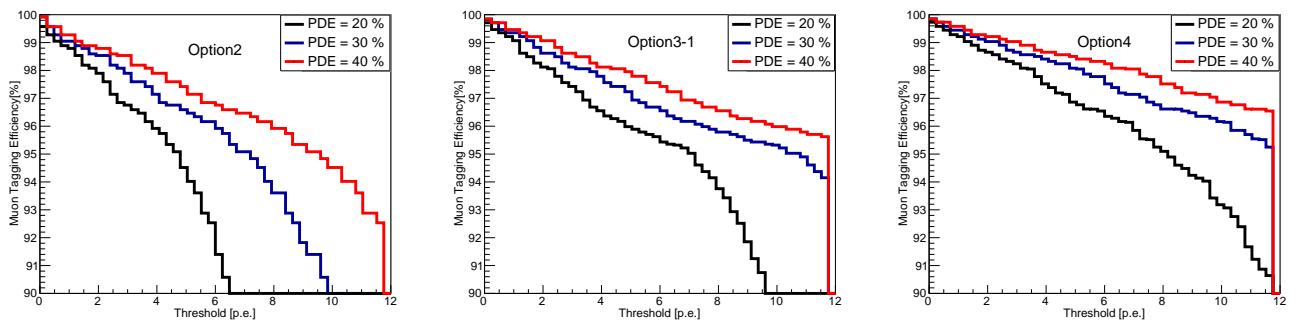


Fig. 11. Relationship between the muon tagging efficiency and the threshold of option 2, option 3-1, option 4, respectively. Where the threshold represents the sum of the SiPM of each PS end over threshold individually. As the threshold increases, the muon tagging efficiency of option 4 decreases the slowest.

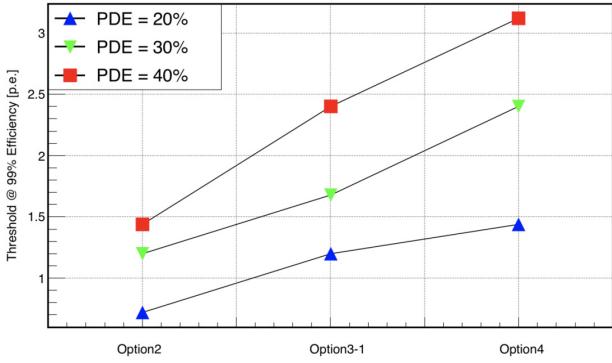


Fig. 12. When the muon tagging efficiency requirement is 99%, the corresponding threshold limits for options 2, 3-1, and 4 respectively. When the PDE of SiPM is 40%, for option 4, the threshold can be set to 3 p.e., which greatly reduces the contribution from dark noise of SiPM.

500 inefficiencies can be indirectly explained by distribution of
501 track length.

502 Figure 14 shows the ratio of the muon events whose track
503 length exceeds certain threshold out of the total events under
504 option 2, option 3-1, and option 4. The trends are almost the
505 same and the difference in ratio was within 1%. The muon
506 generator, the thickness, length, and width of the PS strip are
507 the same for the three configurations, and the declining trend
508 of the event ratio is basically consistent under the three con-
509 figurations. The black horizontal dotted line shows the pro-
510 portion of 99%, from which it can be obtained that the track
511 length threshold of the corresponding black vertical dotted
512 line is 3 mm. There are 1% muon events whose track length
513 in the PS strip is less than 3 mm. Since the thickness of the PS
514 is 20 mm, it is certain that the 1% muon events are incident
515 at a large zenith angle and pass through the edge of the PS,
516 which is often referred to as the edge event. Then, if building
517 a module of PS to stagger the PS up, down, left and right,
518 these edge events should also leave longer tracks on the next
519 layer or the PS next to them, so as to be triggered. Thus, for a
520 PS strip, the edge event may not be tagged(detected), but for
521 the whole veto module system, it can be tagged. Therefore,
522 under the same threshold, the muon tagging efficiency of a
523 single layer module is higher than that of a single PS strip.

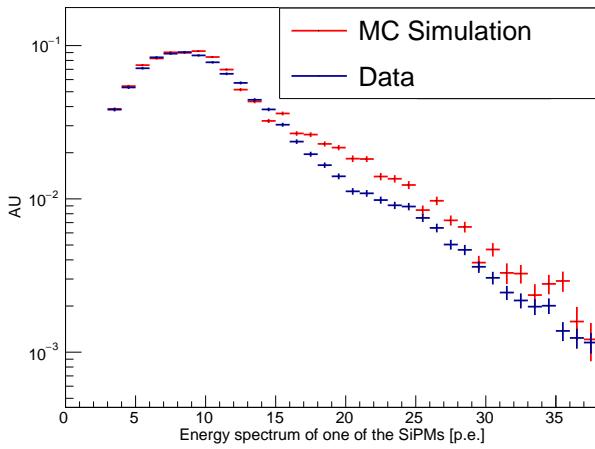


Fig. 13. Comparison between the simulation and experiment in terms of the energy spectrum for option 4. Assuming 40% SiPM PDE in MC. The PDE of the SiPM sample provided by the manufacturer is 46% at the wavelength of 420 nm when the working voltage is 37.2 V. The data is recorded when the CR muon monitor is in the center of PS.

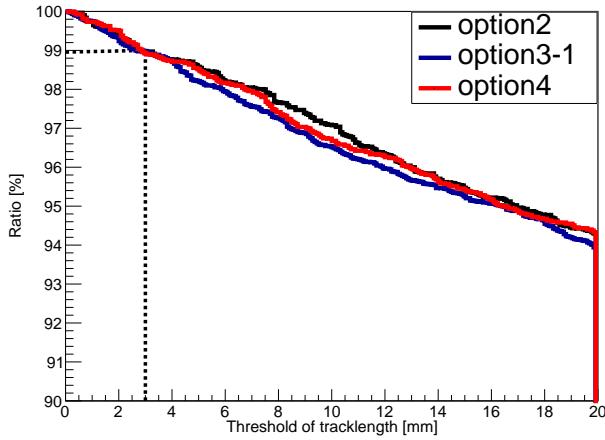


Fig. 14. Event ratio which the track length of incident muon in PS strip exceeds the threshold in the total events under option 2, option 3-1, and option 4, respectively.

Based on this point of view, the JUNO-TAO top veto system was designed, called top veto tracker (TVT).

IV. PERFORMANCE OF JUNO-TAO TVT

Figure 15(a) is the design schema of the TVT. The left panel of Figure 15(a) shows the side view of the whole 4-layer PS, the seam of the upper layer of PS corresponds to the volume of the next layer of PS, which can eliminate the dead

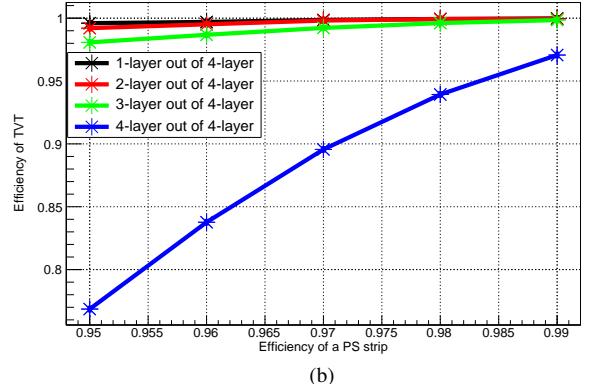
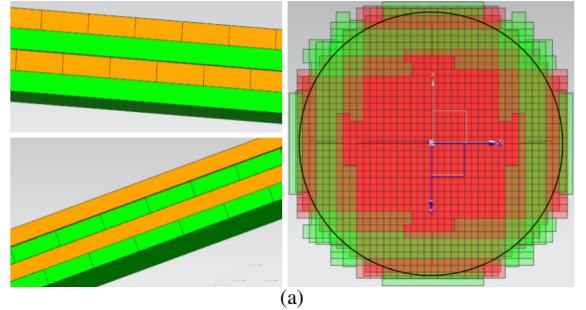


Fig. 15. (a): Design of JUNO-TAO TVT. The left panel is a side view. The right panel is a top view.(b): Muon tagging efficiency of hitting multi-layer PS with that of single.

space. At the same time, it is considered that the track direction of a muon can be determined by this system. Therefore, a 4-layer PS is designed, and the gap between neighboring layers is 2 cm. (It refers to the gap between the SiPM end of one PS and the end of another PS). The right panel of Figure 15(a) shows the top view of the whole 4-layer PS. Different colors represent different layers of PS. Since the central detector of TAO is a ball, under the condition of saving funds and consumables. Our best effort is to cover the ball's projection surface with the area of PS, resulting in the formation of a circle by the four layers.

Figure 15(b) shows the distribution of the muon tagging efficiency of hitting multi-layer PS versus that of a single PS. The different colors represent the efficiencies of hitting the different number of layers of PS. The black line represents the muon tagging efficiency of any one out of four layers. The red represents any two out of four layers. Because the gap between each layer is very small, the red and black lines almost coincide a PS strip efficiency. The blue represents the muon that hits the whole four layers and is also tagged by the four layers of PS. Since each layer has a certain dead space, as the efficiency of the PS strip decreases, the tagging efficiency reduction on all four layers will be most noticeable. The green represents any three out of four layers, from which this conclusion can be obtained: even if the efficiency of a PS strip is as low as 97%, which is equivalent to the threshold of 10 p.e. can be found in Table 3, the muon tagging efficiency

558 of any three out of four layers will still be higher than 99%.
 559 In addition, by simulating the environment background (in-
 560 cluding ^{238}U , ^{232}Th , ^{40}K chain), the background event rate
 561 that can pass the 3 p.e. threshold in any two out of four layers
 562 of PS is 1280 ± 40 Hz and the rate that can pass the 3 p.e.
 563 threshold in any three out of four layers of PS is < 10 Hz.
 564 Therefore, the condition can be selected wherein any three
 565 out of the four layers are triggered to tag muon and reject the
 566 background.

567

V. SUMMARY

568 There are numerous PS with WLS-fiber options with differ-
 569 ent configurations. In this paper, through simulation package,
 570 the differences between different configurations are checked
 571 in detail, for example, for the influence of optical fiber diam-
 572 eter, layout and other factors on transmission performance.
 573 Based on light yield and muon tagging efficiency, the opti-

574 mal configuration of WLS-fiber PS under certain conditions
 575 is obtained. The relevant bench test of option 1 and option 4
 576 has been measured, in terms of the most probable SiPM out-
 577 put, the summed signal at one end of the optimized option 4
 578 is about 32 p.e.. The light yield is 4 times larger than op-
 579 tion 1, which further verified the reliability of the simulation
 580 package.

581 At the same time, inefficiency muon events are also stud-
 582 ied. Finally, the design scheme of JUNO-TAO TVT system
 583 is reached. These conclusions can be given in the simulation:
 584 Assuming 40% SiPM PDE, the muon tagging efficiency of
 585 a single PS strip at a 3 p.e. threshold can still reach 99% in
 586 "AND" mode(i.e., when signals above the threshold were re-
 587 quired from two sides of a PS strip.), the background event
 588 rate that can pass the threshold in any three out of four layers
 589 of PS is less than 10 Hz(Almost zero). When the threshold
 590 is 10 p.e., the muon tagging efficiency of a single PS strip is
 591 97%, The efficiency of 3-layer out of 4-layer of TVT will still
 592 be higher than 99%.

chinaXiv:202302.00256v2

- 593 [1] P.A. Zyla, et al. (Particle Data Group). Review of Particle
 594 Physics. Progress of Theoretical and Experimental Physics. **08**,
 595 08 (2020) [DOI:10.1093/ptep/ptaa104](https://doi.org/10.1093/ptep/ptaa104)
- 596 [2] Mengyun Guan and Ming-Chung Chu and Jun Cao and Kam-
 597 Biu Luk and Changgen Yang. A parametrization of the cosmic-
 598 ray muon flux at sea-level. [DOI:10.48550/arXiv.1509.06176](https://doi.org/10.48550/arXiv.1509.06176)
- 599 [3] Patrignani, C. and others. Particle Data Group. Review
 600 of Particle Physics. Chin. Phys. C. **10**, 100001 (2016).
 601 [DOI:10.1088/1674-1137/40/10/100001](https://doi.org/10.1088/1674-1137/40/10/100001)
- 602 [4] Zi-yi Guo and Lars Bathe-Peters and Shao-min Chen and (JNE
 603 Collaboration). Muon flux measurement at China Jinping Un-
 604 derground Laboratory. Chinese Physics C. **45**, 025001 (2021).
 605 [doi:10.1088/1674-1137/abceae](https://doi.org/10.1088/1674-1137/abceae)
- 606 [5] E. Barbuto and C. Bozza and M. Cozzi and et al. At-
 607 mospheric muon flux measurements at the external site of
 608 the Gran Sasso Lab. Nuclear Instruments and Methods in
 609 Physics Research Section A: Accelerators, Spectrometers,
 610 Detectors and Associated Equipment. **525**, 485-495 (2004).
 611 [doi:https://doi.org/10.1016/j.nima.2004.01.078](https://doi.org/10.1016/j.nima.2004.01.078)
- 612 [6] Trzaska, Wladyslaw Henryk and Slupecki, Maciej and Ban-
 613 dac, et al. Cosmic-ray muon flux at Canfranc Under-
 614 ground Laboratory. European Physical Journal C. **79**, 8 (2019).
 615 [doi:10.1140/epjc/s10052-019-7239-9](https://doi.org/10.1140/epjc/s10052-019-7239-9)
- 616 [7] JUNO Collaboration. JUNO physics and detector. Progress
 617 in Particle and Nuclear Physics. **123**, 0146-6410 (2022).
 618 [doi:https://doi.org/10.1016/j.ppnp.2021.103927](https://doi.org/10.1016/j.ppnp.2021.103927)
- 619 [8] JUNO Collaboration and T. Adam et al. JUNO
 620 Conceptual Design Report. physics.ins-det.
 621 <https://ui.adsabs.harvard.edu/abs/2015arXiv150807166A>
- 622 [9] JUNO Collaboration and Angel Abusleme et al. TAO
 623 Conceptual Design Report: A Precision Mea-
 624 surement of the Reactor Antineutrino Spectrum
 625 with Sub-percent Energy Resolution. physics.ins-det.
 626 <https://ui.adsabs.harvard.edu/abs/2020arXiv200508745J>
- 627 [10] Aprile, E. et al. XENON1T Collaboration. Conceptual de-
 628 sign and simulation of a water Cherenkov muon veto for the
 629 XENON1T experiment. JINST. astro-ph.IM. **9**, P11006 (2014).
 630 [doi:10.1088/1748-0221/9/11/P11006](https://doi.org/10.1088/1748-0221/9/11/P11006)
- 631 [11] Aprile, E. and others. XENON Collaboration. Projected WIMP
 632 sensitivity of the XENONnT dark matter experiment. JCAP.
 633 physics.ins-det. [doi:10.1088/1475-7516/2020/11/031](https://doi.org/10.1088/1475-7516/2020/11/031)
- 634 [12] Christmann, Mirco and others. MAGIX Collaboration. Light
 635 Dark Matter Searches with DarkMESA, PoS, EPS-HEP2021,
 636 129 (2022). [doi:10.22323/1.398.0129](https://doi.org/10.22323/1.398.0129)
- 637 [13] Alexander, T. and others. DarkSide search for dark
 638 matter. JINST. **8**, C11021(2013). [doi:10.1088/1748-0221/8/11/C11021](https://doi.org/10.1088/1748-0221/8/11/C11021)
- 638 [14] Pocar, Andrea. EXO-200, nEXO collaboration. Search-
 639 ing for neutrino-less double beta decay with EXO-200
 640 and nEXO. Nucl. Part. Phys. Proc. **42**, 265-266 (2015)
 641 [doi:10.1016/j.nuclphysbps.2015.06.011](https://doi.org/10.1016/j.nuclphysbps.2015.06.011)
- 642 [15] Tosi, D. EXO collaboration. Search for double beta de-
 643 cay with EXO-200. AIP Conf. Proc. **1560**, 187-189
 644 (2013).[doi:10.1063/1.4826749](https://doi.org/10.1063/1.4826749)
- 645 [16] Gornea, Razvan. EXO-200 collaboration. Double beta de-
 646 cay in liquid xenon. J. Phys. Conf. Ser. **179**, 012004 (2009).
 647 [doi:10.1088/1742-6596/179/1/012004](https://doi.org/10.1088/1742-6596/179/1/012004)
- 648 [17] Birks, John B. The Theory and practice of scintillation count-
 649 ing. (1964). <https://www.slac.stanford.edu/spires/find/books>
- 649 [18] Zhezher, Y. Telescope Array collaboration. Study of Muons
 650 in Ultra-High-Energy Cosmic-Ray Air Showers with the Tele-
 651 scope Array Experiment. Phys. Atom. Nucl. **82**, 685-688
 652 (2020) [doi:10.1134/S1063778819660517](https://doi.org/10.1134/S1063778819660517)
- 653 [19] Erhart, Andreas and others. NUCLEUS collaboration. Devel-
 654 opment of an Organic Plastic Scintillator based Muon Veto Op-
 655 erating at Sub-Kelvin Temperatures for the NUCLEUS Experi-
 656 ment. 19th International Workshop on Low Temperature Detec-
 657 tors. [doi:10.1007/s10909-022-02842-5](https://doi.org/10.1007/s10909-022-02842-5)
- 658 [20] Seo, J. W. and Jeon, E. J. et al. A feasibility study of ex-
 659 truded plastic scintillator embedding WLS fiber for AMoRE-
 660 II muon veto. Nucl. Instrum. Meth.A. **1039**, 167123 (2022).
 661 [doi:10.1016/j.nima.2022.167123](https://doi.org/10.1016/j.nima.2022.167123)
- 662 [21] K.J. Thomas and E.B. Norman and A.R. Smith and Y.D. Chan.
 663 Installation of a muon veto for low background gamma spec-
 664 troscopy at the LBNL low-background facility. Nuclear Instru-
 665 ments and Methods in Physics Research Section A: Ac-
 666 ccelerators, Spectrometers, Detectors and Associated Equipment.
 667 [doi:10.1016/j.nima.2013.05.034](https://doi.org/10.1016/j.nima.2013.05.034)

- chinaXiv:202302.00256v2
- 671 [22] Pla-Dalmau, A. and Bross, A. D. and Mellott, K. L. Low-cost
672 extruded plastic scintillator. *Nucl. Instrum. Meth. A.* **466**, 482-
673 491 (2001). [doi:10.1016/S0168-9002\(01\)00177-2](https://doi.org/10.1016/S0168-9002(01)00177-2)
- 674 [23] Moiseev, A. A. and Hartman, et al. High Efficiency Plastic
675 Scintillator Detector with Wavelength Shifting Fiber Readout
676 for the GLAST Large Area Telescope. *Nucl. Instrum. Meth. A.*
677 **583**, 372-381 (2007). [doi:10.1016/j.nima.2007.09.040](https://doi.org/10.1016/j.nima.2007.09.040)
- 678 [24] Vaishali Manojkumar Thakur and Amit Jain and others. Design
679 and development of a plastic scintillator based whole body be-
680 ta/gamma contamination monitoring system. *Nuclear Science
681 and Techniques.* **32**, 5 (2021). [doi:10.1007/s41365-021-00883-1](https://doi.org/10.1007/s41365-021-00883-1)
- 683 [25] Holm, U. and Wick, K. Radiation Stability of Plastic Scin-
684 tillators and Wave Length Shifters. *IEEE Trans. Nucl. Sci.* **36**,
685 579-583 (1989). [doi:10.1109/23.34504](https://doi.org/10.1109/23.34504)
- 686 [26] Bloise, C. and others. Design, assembly and opera-
687 tion of a Cosmic Ray Tagger based on scintillators and
688 SiPMs. *Nucl. Instrum. Meth. A.* **1045**, 167538 (2023).
689 [doi:10.1016/j.nima.2022.167538](https://doi.org/10.1016/j.nima.2022.167538)
- 690 [27] Buzhan, P. and Karakash, A. Hand-foot monitors for nu-
691 clear plants based on scintillator-WLS-SiPM technology. *J.
692 Phys. Conf. Ser.* **1689**, 012011 (2020). [doi:10.1088/1742-6596/1689/1/012011](https://doi.org/10.1088/1742-6596/1689/1/012011)
- 694 [28] Bugg, W. and Efremenko, Yu. and Vasilyev, S. Large Plas-
695 tic Scintillator Panels with WLS Fiber Readout; Optimiza-
696 tion of Components. *Nucl. Instrum. Meth. A.* **758**, 91-96 (2014).
697 [doi:10.1016/j.nima.2014.05.055](https://doi.org/10.1016/j.nima.2014.05.055)
- 698 [29] Jia-Ning Dong and Yun-Long Zhang and Zhi-Yong Zhang and
699 Dong Liu and Zi-Zong Xu and Xiao-Lian Wang and Shu-
700 Bin Liu. Position-sensitive plastic scintillator detector with
701 WLS-fiber readout. *Nuclear Science and Techniques.* **29**, 117
702 (2018). [doi:10.1007/s41365-018-0449-2](https://doi.org/10.1007/s41365-018-0449-2)
- 703 [30] Y Yang, CP Yang, J Xin, et al. Performance of a plas-
704 tic scintillation fiber dosimeter based on different pho-
705 toelectric devices. *NUCL SCI TECH.* **32**, 120 (2021).
706 <https://doi.org/10.1007/s41365-021-00965-0>
- 707 [31] Adam, T. and others. The OPERA experiment target
708 tracker. *Nucl. Instrum. Meth. A.* **577**, 523-539 (2007).
709 [doi:10.1016/j.nima.2007.04.147](https://doi.org/10.1016/j.nima.2007.04.147)
- 710 [32] Adamson, P. and others. MINOS collaboration. The MINOS
711 scintillator calorimeter system. *IEEE Trans. Nucl. Sci.* **49**, 861-
712 863 (2002). [doi:10.1109/TNS.2002.1039579](https://doi.org/10.1109/TNS.2002.1039579)
- 713 [33] Wang, Ya-Ping and Hou, Chao and Sheng, Xiang-Dong and
714 others. Testing and analysis of the plastic scintillator units
715 for LHAASO-ED. *Rad. Det. Tech. Meth.* **54**, 513-519 (2021).
716 [doi:10.1007/s41605-021-00274-5](https://doi.org/10.1007/s41605-021-00274-5)
- 717 [34] Aharonian, F. and others. LHAASO collaboration. Per-
718 formance test of the electromagnetic particle detectors for the
719 LHAASO experiment. *Nucl. Instrum. Meth. A.* **1001**, 165193
720 (2021). [doi:10.1016/j.nima.2021.165193](https://doi.org/10.1016/j.nima.2021.165193)
- 721 [35] Evans, Justin. MINOS collaboration. The MINOS Experiment:
722 Results and Prospects, *Adv. High Energy Phys.* **2013**, 182537
723 (2013). [doi:10.1155/2013/182537](https://doi.org/10.1155/2013/182537)
- 724 [36] Orsi, Silvio. PAMELA collaboration. PAMELA: A pay-
725 load for antimatter matter exploration and light nuclei as-
726 trophysics. *Nucl. Instrum. Meth. A.* **580**, 880-883 (2007).
727 [doi:10.1016/j.nima.2007.06.051](https://doi.org/10.1016/j.nima.2007.06.051)
- 728 [37] Andreev, V. and others. A high granularity scintillator
729 hadronic-calorimeter with SiPM readout for a linear col-
730 lide detector. *Nucl. Instrum. Meth. A.* **540**, 368-380 (2005).
731 [doi:10.1016/j.nima.2004.12.002](https://doi.org/10.1016/j.nima.2004.12.002)
- 732 [38] Thompson, David J. and Wilson-Hodge, Colleen A. Fermi
733 Gamma-ray Space Telescope. [arXiv:2210.12875](https://arxiv.org/abs/2210.12875). *astro-ph.HE.*
734 [doi:2210.12875](https://arxiv.org/abs/2210.12875)
- 735 [39] S. Procureur. Muon imaging: Principles, technologies and ap-
736 plications. *Nuclear Instruments and Methods in Physics Re-
737 search Section A: Accelerators, Spectrometers, Detectors and
738 Associated Equipment.* [doi:10.1016/j.nima.2017.08.004](https://doi.org/10.1016/j.nima.2017.08.004)
- 739 [40] Morishima, Kunihiro and others. Discovery of a big void in
740 Khufu's Pyramid by observation of cosmic-ray muons. *Nature.*
741 **552**, 7685, 386-390 (2017). [doi:10.1038/nature24647](https://doi.org/10.1038/nature24647)
- 742 [41] Zenoni, Aldo. Historical building stability monitoring by
743 means of a cosmic ray tracking system. *4th International Con-
744 ference on Advancements in Nuclear Instrumentation Mea-
745 surement Methods and their Applications. IEEE Nucl.Sci.
746 Symp.Conf.Rec.* [doi:10.1109/ANIMMA.2015.7465542](https://doi.org/10.1109/ANIMMA.2015.7465542)
- 747 [42] Marteau, J. and Gibert, D. et al. Muons tomography applied
748 to geosciences and volcanology. *Nucl. Instrum. Meth. A.* **695**,
749 23-28 (2012). [doi:10.1016/j.nima.2011.11.061](https://doi.org/10.1016/j.nima.2011.11.061)
- 750 [43] Oguri, S. and Kuroda, Y. et al. Reactor antineutrino mon-
751 itoring with a plastic scintillator array as a new safe-
752 guards method. *Nucl. Instrum. Meth. A.* **757**, 33-39 (2014).
753 [doi:10.1016/j.nima.2014.04.065](https://doi.org/10.1016/j.nima.2014.04.065)
- 754 [44] Georgadze, A. Sh. and Pavlovych, V. M. et al. A re-
755 mote reactor monitoring with plastic scintillation detector.
756 [arXiv:1610.05884](https://arxiv.org/abs/1610.05884). [doi:10.1109/05884/](https://doi.org/10.1109/05884)
- 757 [45] Scovell, P. R. and others. Low background anti-neutrino mon-
758 itoring with an innovative composite solid scintillator detector.
759 2013 IEEE Nuclear Science Symposium and Medical Imag-
760 ing Conference and Workshop on Room-Temperature Semi-
761 conductor Detectors. [doi:10.1109/NSSMIC.2013.682954](https://doi.org/10.1109/NSSMIC.2013.682954)
- 762 [46] Capozzi, Francesco and Lisi, Eligio and Marrone,
763 Antonio. Mapping reactor neutrino spectra from
764 TAO to JUNO. *Phys. Rev. D.* **102**, 056001 (2020).
765 [doi:10.1103/PhysRevD.102.056001](https://doi.org/10.1103/PhysRevD.102.056001)
- 766 [47] S. Agostinelli and others. Geant4—a simulation toolkit. *Nu-
767 clear Instruments and Methods in Physics Research Section A:
768 Accelerators, Spectrometers, Detectors and Associated Equip-
769 ment.* [doi:10.1016/S0168-9002\(03\)01368-8](https://doi.org/10.1016/S0168-9002(03)01368-8)
- 770 [48] Riggi, S. and La Rocca, P. et al. Geant4 simulation of
771 plastic scintillator strips with embedded optical fibers for a
772 prototype of tomographic system. *Nucl. Instrum. Meth. A.*
773 [doi:10.1016/j.nima.2010.10.012](https://doi.org/10.1016/j.nima.2010.10.012)
- 774 [49] Wenzhen XU and Yanfen LIU and Zongquan TAN and Ran
775 XIAO and Wei KONG and Bangjiao YE. Geant4 simulation of
776 plastic scintillators for a prototype uSR spectrometer. *Nuclear
777 Science and Techniques.* **24**, 4 (2013) [doi:10.13538/j.1001-8042/nst.2013.04.011](https://doi.org/10.13538/j.1001-8042/nst.2013.04.011)
- 778 [50] Lecoq, P. Scintillation Detectors for Charged Particles and Phot-
779 tons. *Particle Physics Reference Library.* Springer. Cham, 45-
780 89 (2020). [doi:10.1007/978-3-030-35318-6-3](https://doi.org/10.1007/978-3-030-35318-6-3)
- 781 [51] Min Li and Zhi Min Wang and Cai Mei Liu and Pei Zhi Lu
782 and Guang Luo and Yuen Keung Hor and Jin Chang Liu and
783 Chang-Gen Yang. Performance of compact plastic scintillator
784 strips with wavelength shifting fibers using a photomultiplier
785 tube or silicon photomultiplier readout. *Nuclear Science and
786 Techniques.* **34**, 2 (2023). [doi:10.1007/s41365-023-01175-6](https://doi.org/10.1007/s41365-023-01175-6)
- 787 [52] Yang, Hang and Luo, Guang. et al. MuGrid: A scin-
788 tillator detector towards cosmic muon absorption imag-
789 ing. *Nucl. Instrum. Meth. A.* **1042**, 167402 (2022).
790 [doi:10.1016/j.nima.2022.167402](https://doi.org/10.1016/j.nima.2022.167402)
- 791 [53] Hoton Technology Co. Beijing Hoton Nuclear Technology Co.,
792 Ltd. [doi://www.hoton.com.cn](https://www.hoton.com.cn)
- 793 [54] Tur, Clarisse and Solov'yev, Vladimir and Flamanc, Jeremy.
794 Temperature characterization of scintillation detectors us-
795 ing solid-state photomultipliers for radiation monitoring ap-
796

- 797 plications. *Nucl. Instrum. Meth. A.* **620**, 351-358 (2010).
798 [doi:10.1016/j.nima.2010.03.141](https://doi.org/10.1016/j.nima.2010.03.141)
- 799 [55] Dietz Laursonn, Erik. Detailed Studies of Light Transport
800 in Optical Components of Particle Detectors. Aachen, Tech.
801 Hochsch. doi:inspirehep.net/literature/1505685
- 802 [56] Qian, Xiang-Li and Sun, Hui-Ying and Liu, Cheng and Wang,
803 Xu and Martineau-Huynh, Olivier. Simulation study on perfor-
804 mance optimization of a prototype scintillation detector for the
805 GRANDProto35 experiment. *Nucl. Sci. Tech.* **32**, 51 (2021).
806 [doi:10.1007/s41365-021-00882-2](https://doi.org/10.1007/s41365-021-00882-2)
- 807 [57] Gaisser, Thomas. Cosmic-Ray Showers Reveal Muon Mystery.
808 *APS Physics*. **9**, 125 (2016). [doi:10.1103/Physics.9.125](https://doi.org/10.1103/Physics.9.125)
- 809 [58] Shukla, Prashant and Sankrith, Sundaresh. Energy
810 and angular distributions of atmospheric muons at
811 the Earth. *Int. J. Mod. Phys. A.* **33**, 1850175 (2018).
812 [doi:10.1142/S0217751X18501750](https://doi.org/10.1142/S0217751X18501750)
- 813 [59] Semiconductor Components Industries, LLC. Cherry Semicon-
814 ductor. (1999-2023). [doi:onsemi.com](https://doi.onsemi.com)